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SPACE SYSTEM ASSOCIATES

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(NASA-CR-144780) LONG PERIOD COUPLING TERMS  
FOR LAGRANGE'S EQUATIONS (Space System  
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For Lagrange's Equations "

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# LONG PERIOD COUPLING TERMS IN LAGRANGE'S EQUATIONS

## Introduction

In this report the long period terms arising from the short-short period coupling of zonal harmonics are derived for Lagrange's Equations. The formulation is general so that the results are valid for any pairs of zonal harmonics  $J_n$  where  $n$  and  $\ell$  are arbitrary.

Formulas are given to generate the various functions and integrals needed for the results given in this report. Checks have been made against the work of Kozai, Reference 1.

This paper is a generalization of that portion of the work of Berger, Reference 2, which deals with the long period coupling effect of certain pairs of zonal harmonics.

## Analysis

### Lagrange's Equations

The equations of Lagrange are:

$$\frac{da}{dt} = \frac{2a^{\frac{1}{2}}}{\mu^{\frac{1}{2}}} \frac{\partial R}{\partial M}$$

$$\frac{de}{dt} = \frac{1-e^2}{\mu^{\frac{1}{2}} a^{\frac{1}{2}}} \frac{\partial R}{\partial M} - \frac{\sqrt{1-e^2}}{\mu^{\frac{1}{2}} a^{\frac{1}{2}}} \frac{\partial R}{\partial w} \quad (1)$$

$$\frac{di}{dt} = \frac{\cos i}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} \sqrt{1-e^2}} \frac{\partial R}{\partial w}$$



$$\frac{d\omega}{dt} = - \frac{\cos i}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} \sqrt{1-e^2} \sin i} \frac{\partial R}{\partial i} + \frac{\sqrt{1-e^2}}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} e} \frac{\partial R}{\partial e}$$

$$\frac{d\Omega}{dt} = \frac{1}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} \sqrt{1-e^2} \sin i} \frac{\partial R}{\partial i}$$

$$\frac{dM}{dt} = \mu^{\frac{1}{2}} a^{\frac{3}{2}} - \frac{1-e^2}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} e} \frac{\partial R}{\partial e} - \frac{2a^{\frac{1}{2}}}{\mu^{\frac{1}{2}}} \frac{\partial R}{\partial a}$$

The Disturbing Function

The short period disturbing function is given by

$$R = \sum_{n=2}^N R_n$$

where  $\frac{n-ODN}{2}$

$$R_n = \sum_{q=0}^{\frac{n-ODN}{2}} \frac{\mu J_2}{a^{n+1}} B_{n,q} \left[ \left( \frac{a}{r} \right)^{n+1} \cos(\alpha + qf) - a_{n,q} \cos \alpha \right] \quad (2)$$

$$q = 2i + ODN$$

where  $ODN = 0$  for  $n$  even  
 $= 1$  for  $n$  odd

$$\alpha = q\omega - ODN \frac{\pi}{2}$$

This notation is introduced to keep the form of the disturbing function invariant. This we do not need to write separate equations for  $R_n$  for the cases  $n$  even or odd.

Since we deal with coupling of zonals  $n$  and  $l$  the function  $R_2$  is defined similarly to  $R_n$  where  $n, q, i, ODN$  and  $\alpha$  are replaced by  $l, s, j, ODL$  and  $\beta$ , respectively.

The quantities  $B_{n,q}$  and  $a_{n,q}$  are the inclination and averaging functions respectively and are given in Appendix 1.

The product  $a_{n,q} \cos \alpha$  in Equation 1, is the average of the term  $\left(\frac{a}{r}\right)^{n+1} \cos(\alpha + q f)$  taken over the mean anomaly.

#### Expansions For Lagrange's Equations

The Taylor series expansion for a typical Keplerian element  $E_i$  needed to find the coupling terms in Equations (1) is

$$\frac{dE_i}{dt} = \sum_{n=2}^N \left( \frac{f_i}{\delta i} \right) + \left[ \sum_{n=2}^N \sum_{l=2}^L \left\{ \left( a \frac{\partial f_i}{\partial a} \right) \left( \frac{\delta a}{a} \right)_l + \left( \frac{\partial f_i}{\partial e} \right) (\delta e)_l + \left( \frac{\partial f_i}{\partial i} \right) (\delta i)_l + \left( \frac{\partial f_i}{\partial \omega} \right) (\delta \omega)_l + \left( \frac{\partial f_i}{\partial M} \right) (\delta M)_l \right\} \right] \quad (3)$$

$f_i$  represents the functions appearing on the right hand side of Equations (1) for the  $i^{\text{th}}$  element. For example for the inclination

$$f_i = \frac{\cos i}{a^{1/2} \eta \sin i} \frac{\partial R_n}{\partial \omega}$$

The symbol  $\delta$  attached to an element indicates the perturbations in that element. The subscripts  $n$  and  $l$  pertain to the zonal harmonics considered.

General formulas for  $f_i$ ,  $\frac{\partial f_i}{\partial \epsilon_j}$ , and for  $\delta \epsilon_i$  are given in appendices 3 and 4.

In Equation (3),  $(N-1)(L-1)$  coupling terms arise. Formulas for the general  $J_n - J_2$  long period term are derived by averaging the quantities enclosed by the brackets in Equation (3), using the functions defined by the appendices. The results are functions of  $\alpha$  and  $\beta$  which are defined in connection with Equation (2).

It should be noted that long period perturbations due to  $J_n - J_2$  coupling do not occur in the semi major axis. This was proved by Kozai, Reference 4. His proof while given only for  $J_2 - J_2$  coupling can be readily extended for 2<sup>nd</sup> order coupling of the form  $J_n - J_2$ . Long period terms in the semi major axis result from higher order coupling of the form  $J_n - J_2 - J_m$ , which is beyond the scope of this report. However for the special but important case for long period terms of order  $J_2^2$  the results of Berger, Reference 5, are available.

#### Evaluation of the Expansions For Lagrange's Equations

The mean values of the functions enclosed by the brackets in Equation (3) are now given for each of the Keplerian elements under consideration. Since certain types of trigonometric products recur frequently they are given in Appendix 5.

In general the functions to be averaged are of the form

$$\left(\frac{a}{r}\right)^{p+1} \cos(\alpha + q f)$$

so that the average values are readily by means of the averaging function,  $a_{p-1, q}$  described in Appendix 1.

When the subscripts of the averaging function are functions of a parameter e.g.  $a_{v(t), \phi(t)}$ , then it is useful to find the maximum value of  $t$ ,  $t_{max}$ , which will yield non-zero terms for the averaging function. Then  $t_{max}$  is found from the relation

$$v(t) - |\phi(t)| = 0$$

For example if  $v(t) = n$ ,  $\phi(t) = q - 1 - t$ , then

$$t_{max} = n + q - 1$$

In some of the series a term occurs in which  $s - k$  is a divisor, where  $k$  is a summation index. These series have been derived so that the correct sum is obtained by omitting the term for which  $s - k = 0$ .

#### The Inclination Equation

The expansion for the inclination equation corresponding to Equation (3) is given by

$$\begin{aligned} \frac{di}{dt} = & \left( f_3 \frac{\partial R}{\partial \omega} \right) + \left[ f_3 \left\{ \left( a \frac{\partial^2 R}{\partial \omega \partial a} \right)_n \left( \frac{\partial a}{a} \right)_2 + \left( \frac{\partial^2 R}{\partial \omega \partial e} \right)_n (\partial e)_2 + \left( \frac{\partial^2 R}{\partial \omega \partial i} \right)_n (\partial i)_2 \right. \right. \\ & \left. \left. + \left( \frac{\partial^2 R}{\partial \omega^2} \right)_n (\partial \omega)_2 \right\} + \left( \frac{\partial R}{\partial \omega} \right)_n \left\{ a \frac{\partial f_3}{\partial a} \left( \frac{\partial a}{a} \right)_2 + \frac{\partial f_3}{\partial e} (\partial e)_2 + \frac{\partial f_3}{\partial i} (\partial i)_2 \right\} \right] \end{aligned}$$

$$f_3 = \frac{\cos i}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} \eta \sin i}$$

$$a \frac{\partial f_3}{\partial a} = - \frac{\cos i}{2 \mu^{\frac{1}{2}} a^{\frac{1}{2}} \eta \sin i}$$

$$\frac{\partial f_3}{\partial e} = \frac{e \cos i}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} \eta^3 \sin i}$$

$$\frac{\partial f_3}{\partial i} = - \frac{1}{\mu^{\frac{1}{2}} a^{\frac{1}{2}} \eta \sin^2 i}$$

$$\eta = \sqrt{1 - e^2}$$

Using the functions given in the appendices,

$$1. \left( a \frac{\partial R}{\partial \omega \partial a} \frac{\partial a}{\partial a} \right) = \frac{\mu J_n J_e}{a^{n+L+1}} \sum_q \sum_s 2(n+1) g B_{l,s} B_{n,q} (\overline{A_n} A_e - \overline{A_n'} \overline{A_e'})$$

The bars over the functions indicate averages over the mean anomaly and consequently define the long period terms.

$\sum_q, \sum_s$  are abbreviations for  $\sum_{q=0}^{\frac{n-ODN}{2}}, \sum_{s=0}^{\frac{l-ODL}{2}}$  respectively, where we recall that

$$q = 2i + ODN \quad l = 0, 1, 2, \dots, \frac{n-ODN}{2}$$

$$s = 2j + ODL \quad j = 0, 1, 2, \dots, \frac{l-ODL}{2}$$

$$\alpha = q\omega - ODN \frac{\pi}{2}$$

$$\beta = s\omega - ODL \frac{\pi}{2}$$

The functions of the type  $\overline{A_n}, A_n$  etc. are defined in appendix 2.



$$\overline{A'_n A'_l} - \overline{A'_n A'_e} = \frac{1}{2} \left\{ \sin(\alpha + \beta) [a_{n+l, q+1} - a_{n+1, q} a_{l-1, 0}] \right. \\ \left. + \sin(\alpha - \beta) [a_{n+l, q-1} - a_{n+1, q} a_{l-1, 0}] \right\}$$

$$2. \frac{\partial^2 R}{\partial \omega \partial e} (\delta e)_i = - \frac{\mu \eta^2 J_2 J_e}{e a^{n+1}} \sum_q \sum_s q B_{n, q} B_{l, 0} [\overline{F'_n A_s} - \overline{F'_n A_e}]$$

$$\overline{F'_n A_e} = \sum_{d=1}^4 \frac{c(d)}{2} \left\{ \sin(\alpha + \beta) (a_{v(d), \phi(d)}) + \sin(\alpha - \beta) (a_{v(d), \rho(d)}) \right\}$$

The parameters for  $\overline{F'_n A_e}$  are given in Table 1.

TABLE 1 (Parameters for  $\overline{F'_n A_e}$ )

d	c(d)	v(d)	$\phi(d)$	$\rho(d)$
1	$\frac{n+1+q}{2}$	$n+l+1$	$q+s+1$	$q-s+1$
2	$\frac{q}{2\eta^2}$	$n+l$	$q+s+1$	$q-s+1$
3	$\frac{n+1-q}{2}$	$n+l+1$	$q+s-1$	$q-s-1$
4	$-\frac{q}{2\eta^2}$	$n+l$	$q+s-1$	$q-s-1$

$$\overline{F'_n A_e} = \sin \alpha \cos \beta \left\{ b_{n+1, q} - \frac{(2n-1)}{\eta} a_{n+1, q} \right\} a_{l-1, 0}$$



$$(3) \frac{\delta R}{\delta \omega \delta e} (\delta e)_2 = \frac{\mu \eta \lambda \lambda_e}{e \lambda^{n+1}} \sum_q \sum_s s a_{s,q} B_{n,q} B_{e,s} (F'_n B_{kp}) \sin \beta$$

$$\overline{F'_n B_{kp}} = \sum_{d=1}^4 c(d) \left\{ \sum_{t=1}^{t_1} \gamma_t a_{v(d), \phi(d)} - \sum_{t=1}^{t_2} \gamma_t a_{v(d), \tau(d)} \right\}$$

The parameters for  $\overline{F'_n B_{kp}}$  are given in Table 2.

TABLE 2 (Parameters for  $\overline{F'_n B_{kp}}$ )  
( $c(d)$  same as in Table 1)

$d$	$v(d)$	$\phi(d)$	$\tau(d)$	$t_1$	$t_2$
1	$n$	$q+1+t$	$q+1-t$	$n-q-1$	$n+q+1$
2	$n-1$	$q+1+t$	$q+1-t$	$n-q-2$	$n+q$
3	$n$	$q-1+t$	$q-1-t$	$n-q+1$	$n+q-1$
4	$n-1$	$q-1+t$	$q-1-t$	$n-q$	$n+q-2$

$$(4) \frac{\delta R}{\delta \omega \delta e} (\delta e)_3 = \frac{\mu \eta \lambda \lambda_e}{e \lambda^{n+1}} \sum_q \sum_s \sum_{k=0}^{k=1} e_k s a_{s,q,k} B_{n,q} B_{e,s} (\overline{F'_n C_e} - \overline{F'_n \bar{C}_e})$$

$$\overline{F'_n C_e} = \sum_{d=1}^4 \frac{c(d)}{2} \left\{ a_{v(d), \phi(d)} \frac{\sin(\alpha+\beta)}{s+k} + a_{v(d), \tau(d)} \frac{\sin(\alpha-\beta)}{s+k} \right. \\ \left. + a_{v(d), \sigma(d)} \frac{\sin(\alpha+\beta)}{s-k} + a_{v(d), \psi(d)} \frac{\sin(\alpha-\beta)}{s-k} \right\}$$

The parameters for  $\overline{F'_n C_e}$  are given in Table 3.

TABLE 3 (Parameters for  $\overline{F'_n C_e}$ )  
( $c(d)$  same as in Table 1)

$d$	$v(d)$	$\phi(d)$	$\tau(d)$	$\sigma(d)$	$\psi(d)$
1	$n$	$q+1+s+k$	$q+1-s-k$	$q+1+s-k$	$q+1-s+k$
2	$n+1$	$q+1+s+k$	$q+1-s-k$	$q+1+s-k$	$q+1-s+k$
3	$n$	$q-1+s+k$	$q-1-s-k$	$q-1+s-k$	$q-1-s+k$
4	$n-1$	$q-1+s+k$	$q-1-s-k$	$q-1+s-k$	$q-1-s+k$

$$\overline{F}_n \overline{C}_e = \sin \alpha \cos \beta (\gamma_{s+k} + \gamma_{s-k})$$

$$5. \overline{\frac{\partial^2 R}{\partial s \partial i}} (\delta i)_1 = - \frac{\mu J_n J_e \cos i}{\eta a^{n+1/2} \sin i} \sum_q \sum_s q^2 s a_{e-1,s} B_{n,q} B_{2,s} \cos \beta (\overline{A}_n \overline{B}_{kp})$$

$$\overline{A}_n \overline{B}_{kp} = \cos \alpha \left\{ \sum_{t=1}^{n-q+1} \gamma_t a_{n-1,q+t} - \sum_{t=1}^{n-q-1} \gamma_t a_{n-1,q-t} \right\}$$

$$6. \overline{\frac{\partial^2 R}{\partial \omega \partial e}} (\delta e)_2 = - \frac{\mu J_n J_e \cos i}{\eta a^{n+1/2} \sin i} \sum_q \sum_s \sum_{k=0}^{e-1} e_k q^2 s a_{e-1,k} B_{n,q} C_{2,s} (\overline{A}_n \overline{C}_2 - \overline{A}_n \overline{C}_2)$$

$$\overline{A}_n \overline{C}_2 = \frac{1}{2} \left\{ \frac{\sin(\alpha+\beta)}{s+k} a_{n-1,q+s+k} + \frac{\sin(\alpha-\beta)}{s+k} a_{n-1,q-s+k} + \frac{\sin(\alpha+\beta)}{s-k} a_{n-1,q+s-k} \right. \\ \left. + \frac{\sin(\alpha-\beta)}{s-k} a_{n-1,q-s-k} \right\}$$

$$\overline{A}_n \overline{C}_2 = \sin \alpha \cos \beta a_{n-1,q} a_{e-1,k} (\gamma_{s+k} + \gamma_{s-k})$$

$$7. \overline{\frac{\partial^2 R}{\partial \omega^2}} (\delta \omega)_1 = - \frac{\mu J_n J_e}{\eta^2 a^{n+1/2}} \sum_q \sum_s q^2 B_{n,q} W \cos \beta (\overline{A}_n \overline{B}_{kp})$$

$$\overline{A}_n \overline{B}_{kp} = - \sin \alpha \left\{ \sum_{t=1}^{n-q+1} \gamma_t a_{n-1,q+t} - \sum_{t=1}^{n-q-1} \gamma_t a_{n-1,q-t} \right\}$$

$$8. \overline{\frac{\partial^2 R}{\partial \omega^2}} (\delta \omega)_2 = - \frac{\mu J_n J_e}{\eta^2 a^{n+1/2}} \sum_q \sum_s \sum_{k=0}^{e-1} e_k q^2 B_{n,q} W (\overline{A}_n \overline{C}'_2 - \overline{A}_n \overline{C}_2)$$

$$\overline{A}_n \overline{C}'_2 = \frac{1}{2} \left\{ \frac{\sin(\alpha+\beta)}{s+k} a_{n-1,q+s+k} - \frac{\sin(\alpha-\beta)}{s+k} a_{n-1,q-s+k} + \frac{\sin(\alpha+\beta)}{s-k} a_{n-1,q+s-k} \right. \\ \left. - \frac{\sin(\alpha-\beta)}{s-k} a_{n-1,q-s-k} \right\}$$

$$\overline{A}_n \overline{C}'_2 = \sin \beta \cos \alpha (\gamma_{s+k} + \gamma_{s-k})$$

$$9. \frac{\partial^2 R}{\partial \omega^2}(\delta \omega)_3 = - \frac{\mu J_n J_p}{2 \varepsilon a^{n+l+1}} \sum_q \sum_s q^2 B_{n,q} B_{l,s} (\overline{A_n D'_{l,s}} - \overline{A_n D'_{l,s}})$$

$$\overline{A_n D'_{l,s}} = \frac{\sin(\alpha + \beta)}{2} [a_{n+l-1, s+1+q} - a_{n+l-1, s-1+q}]$$

$$+ \frac{\sin(\beta - \alpha)}{2} [a_{n+l-1, s+1-q} - a_{n+l-1, s-1-q}]$$

$$\overline{A_n D'_{l,s}} = \sin \beta \cos \alpha [a_{l-2, s+1} - a_{l-2, s-1}] a_{n+1, q}$$

$$10. \frac{\partial^2 R}{\partial \omega^2}(\delta \omega)_4 = - \frac{\mu J_n J_p}{2 \varepsilon \eta^2 a^{n+l+1}} \sum_q \sum_s q^2 B_{n,q} B_{l,s} (\overline{A_n D'_{l,s}} - \overline{A_n D'_{l,s}})$$

$$\overline{A_n D'_{l,s}} = \frac{\sin(\beta + \alpha)}{2} [a_{n+l-2, s+1+q} - a_{n+l-2, s-1+q}]$$

$$+ \frac{\sin(\beta - \alpha)}{2} [a_{n+l-2, s+1-q} - a_{n+l-2, s-1-q}]$$

$$\overline{A_n D'_{l,s}} = \sin \beta \cos \alpha [a_{l-3, s+1} - a_{l-3, s-1}] a_{n+1, q} \text{ for } l \geq 3$$

$$\overline{A_n D'_{l,s}} = \frac{\sin \beta \cos \alpha}{2} (a_{n+1, q}) [\overline{\cos(s+1)} - \overline{\cos(s-1)}]$$

$$+ \frac{\eta^2}{2} \{ \overline{\cos(s+2)} - \overline{\cos(s-2)} \} \text{ for } l=2$$

$$11. \frac{\partial^2 R}{\partial \omega \partial \eta}(\delta \eta)_1 = - \frac{\mu J_n J_p}{\varepsilon \eta^{l+3} a^{n+l+1}} \sum_q \sum_s q^2 B'_{l-1, s} B_{n,q} B_{kp} \cos \beta (\overline{E_n B_{kp}})$$

$$\overline{E_n B_{kp}} = \sum_{d=1}^3 c(d) \left\{ \sum_{t=1}^{t_1} x_t a_{v(d), \phi(d)} - \sum_{t=1}^{t_2} x_t a_{v(d), \phi(d)} \right\} \text{ and}$$

The parameters for  $\overline{E_n B_{kp}}$  are given in Table 4.

TABLE 4 (Parameters for  $\overline{E_n B_{kp}}$ )

$d$	$c(d)$	$v(d)$	$\phi(d)$	$t_1$	$\phi(d)$	$t_2$
1	$-\frac{e}{\eta} \left( \frac{n+1}{2} \right)$	$n$	$q+1-t$	$n+q+1$	$q+1+t$	$n-q-1$
2	$\frac{e}{\eta} \left( \frac{n+1}{2} \right)$	$n$	$q-1-t$	$n+q-1$	$q-1+t$	$n-q+1$
3	$-\frac{q}{\eta}$	$n+1$	$q-t$	$n+q+1$	$q+t$	$n-q+1$

$$12. \frac{\partial^2 R}{\partial \omega \partial H} (\delta H)_2 = - \frac{\mu J_n J_e}{e \gamma^{2l+3} a^{2l+1}} \sum_q \sum_s \sum_{k=0}^{E-1} \epsilon_k q b'_{k,l,k} B_{n,q} B_{2,s} (\overline{E_n C'_2})$$

$$\overline{E_n C'_2} = \sum_{d=1}^3 \frac{c(d)}{2} \left\{ \frac{\sin(\beta+d)}{s+k} a_v(d), p(d) + \frac{\sin(\beta-d)}{s+k} a_v(d), \sigma(d) \right. \\ \left. + \frac{\sin(\beta+d)}{s-k} a_v(d), p(d) + \frac{\sin(\beta-d)}{s-k} a_v(d), \sigma(d) \right\}$$

The parameters for  $\overline{E_n C'_2}$  are given in Table 5.

TABLE 5 ( Parameters for  $\overline{E_n C'_2}$ ,  $c(d)$  same as Table 4)

$d$	$\phi(d)$	$p(d)$	$\sigma(d)$	$\tau(d)$
1	$s+k+q+1$	$s-k+q+1$	$s+k-q-1$	$s-k-q-1$
2	$s+k+q-1$	$s-k+q-1$	$s+k-q+1$	$s-k-q+1$
3	$s+k+q$	$s-k+q$	$s+k-q$	$s-k-q$

$$13. \frac{\partial^2 R}{\partial \omega \partial H} (\delta H)_3 = - \frac{\mu J_n J_e \gamma}{e a^{n+l+1}} \sum_q \sum_s q B_{n,q} B_{2,s} (\overline{E_n D'_{12}})$$

$$\overline{E_n D'_{12}} = \sum_{d=1}^3 \frac{c(d)}{4} \left\{ \sin(\beta+d) [a_v(d), \tau(d) - a_v(d), p(d)] \right. \\ \left. + \sin(\beta-d) [a_v(d), \sigma(d) - a_v(d), \tau(d)] \right\}$$

The parameters for  $\overline{E_n D'_{12}}$  are given in Table 6.

TABLE 6 ( Parameters for  $\overline{E_n D'_{12}}$ ,  $c(d)$  same as Table 4)

$d$	$\gamma(d)$	$\phi(d)$	$p(d)$	$\sigma(d)$	$\tau(d)$
1	$n+l$	$s+q+2$	$s+q$	$s-q$	$s-q+2$
2	$n+l$	$s+q$	$s+q-2$	$s-q-2$	$s-q$
3	$n+l-1$	$s+q+1$	$s+q-1$	$s-q+1$	$s-q-1$



$$14. \frac{\partial^2 R}{\partial \omega \partial \eta} (\overline{M})_y = -\frac{\mu J_n J_e}{e \eta a^{n+l+1}} \sum_q \sum_s q B_{n,q} B_{l,s} (\overline{E_n D'_{le}})$$

$$\overline{E_n D'_{le}} = \sum_{d=1}^3 \frac{c(d)}{4} \left\{ \sin(\beta + \alpha) [a_{v(d), \phi(d)} - a_{v(d), \tau(d)}] \right. \\ \left. + \sin(\beta - \alpha) [a_{v(d), \sigma(d)} - a_{v(d), \tau(d)}] \right\}$$

The parameters  $c(d)$ ,  $\phi(d)$ ,  $\sigma(d)$ ,  $\tau(d)$ ,  $\alpha(d)$ , are the same as in Table 5. The parameter  $v(d)$  is given in Table 6.

TABLE 6 ( Parameters for  $\overline{E_n D'_{le}}$  )

d	v(d)
1	n+l-1
2	n+l-1
3	n+l-2

$$15. \frac{\partial R}{\partial \omega} \frac{\partial a}{\partial \eta} = \frac{-2\mu J_n J_e}{a^{n+l+1}} \sum_q \sum_s q B_{n,q} B_{l,s} (\overline{A'_n A_e} - \overline{A'_n \overline{A_e}})$$

$$\overline{A'_n A_e} - \overline{A'_n \overline{A_e}} = \frac{\sin(\alpha + \beta)}{2} [a_{n+l, q+s} - a_{l, s} a_{n, q}] \\ + \frac{\sin(\alpha - \beta)}{2} [a_{n+l, q-s} - a_{l, s} a_{n, q}]$$

$$16. \frac{\partial R}{\partial \omega} (\overline{S_e})_l = \frac{-\eta \mu J_n J_e}{e a^{n+l+1}} \sum_q \sum_s q B_{n,q} B_{l,s} (\overline{A'_n A_e} - \overline{A'_n \overline{A_e}})$$

The quantity  $\overline{A'_n A_e} - \overline{A'_n \overline{A_e}}$  is defined in Equation (15) above.

$$17. \frac{\partial R}{\partial \omega} (\overline{S_e})_v = \frac{-\eta \mu J_n J_e}{e a^{n+l+1}} \sum_q \sum_s q s a_{l, s} B_{n,q} B_{l,s} \sin \beta (\overline{A'_n B_{vf}})$$

The quantity  $\overline{A'_n B_{vf}}$  is defined in Equation (5) above.

$$18. \quad \overline{\frac{\partial R}{\partial \omega}(\delta e)_2} = \frac{\eta \mu J_n J_e}{e a^{n+e+1}} \sum_q \sum_s \sum_{k=0}^{e-1} \epsilon_k q s a_{e-1,k} B_{n,q} B_{e,s} (\overline{A'_n C_e} - \overline{A_n C'_e})$$

The quantity  $\overline{A'_n C_e} - \overline{A_n C'_e}$  is defined in Equation (6) above.

$$19. \quad \overline{\frac{\partial R}{\partial \omega}(\delta i)_1} = \frac{\mu J_n J_e \cos i}{\eta a^{n+e+1} \sin i} \sum_q \sum_s q s a_{e-1,s} B_{n,q} B_{e,s} \sin \beta (\overline{A'_n B_{e\beta}})$$

The quantity  $\overline{A'_n B_{e\beta}}$  is defined in Equation (5) above.

$$20. \quad \overline{\frac{\partial R}{\partial \omega}(\delta i)_2} = \frac{-\mu J_n J_e \cos i}{\eta a^{n+e+1} \sin i} \sum_q \sum_s \sum_{k=0}^{e-1} \epsilon_k q s a_{e-1,k} B_{n,q} B_{e,s} (\overline{A'_n C_e} - \overline{A_n C'_e})$$

The quantity  $\overline{A'_n C_e} - \overline{A_n C'_e}$  is defined in Equation (6) above.

The long period terms for the inclination equation are the sum of the 20 equations given above.

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## The Eccentricity Equation

The expansion for the eccentricity equation is given by the expression,

$$\left[ \frac{de}{dt} \right] = \left[ f_{21} \left\{ \left[ \frac{\partial R}{\partial M} \right] \right\} + \frac{\partial R}{\partial M} \left\{ a \frac{\partial f_{11}}{\partial a} \frac{\delta a}{a} + \frac{\partial f_{11}}{\partial e} \delta e \right\} - f_{22} \left\{ \left[ \frac{\partial R}{\partial \omega} \right] \right\} - \frac{\partial R}{\partial \omega} \left\{ a \frac{\partial f_{22}}{\partial a} \frac{\delta a}{a} + \frac{\partial f_{22}}{\partial e} \delta e \right\} \right]$$

Where  $\left[ \frac{\partial R}{\partial M} \right]$  represents the expansion,

$$\left[ \frac{\partial R}{\partial M} \right] = a \frac{\partial^2 R}{\partial M \partial a} \frac{\delta a}{a} + \frac{\partial^2 R}{\partial M \partial e} \delta e + \dots + \frac{\partial^2 R}{\partial M^2} \delta M, \text{ and similarly}$$

$$\left[ \frac{\partial R}{\partial \omega} \right] = a \frac{\partial^2 R}{\partial \omega \partial a} \frac{\delta a}{a} + \frac{\partial^2 R}{\partial \omega \partial e} \delta e + \dots + \frac{\partial^2 R}{\partial \omega \partial M} \delta M.$$

$$f_{21} = \frac{\eta^2}{\mu^2 a^2 e}$$

$$a \frac{\partial f_{21}}{\partial a} = \frac{-\eta^2}{2 \mu^2 a^2 e}$$

$$\frac{\partial f_{21}}{\partial e} = -\frac{(1+e^2)}{\mu^2 a^2 e}$$

$$f_{22} = \frac{-\eta}{\mu^2 a^2 e}$$

$$a \frac{\partial f_{22}}{\partial a} = \frac{\eta}{2 \mu^2 a^2 e}$$

$$\frac{\partial f_{22}}{\partial e} = -\frac{1}{\mu^2 a^2 e^2 \eta}$$

It can be shown by an extension of the method of Kozai, reference 4, that no long period coupling terms arise from  $\left[\frac{\partial R}{\partial M}\right]$  and the product  $\frac{\partial a}{\partial M} \frac{\partial R}{\partial M}$ . Below we give the long period terms for the products  $\frac{\partial R}{\partial e}(\delta e)_1$  and  $\frac{\partial R}{\partial q}(\delta e)_2$  since the remaining products of  $\frac{de}{dt}$  have been given in the inclination equation.

$$1. \quad \overline{\frac{\partial R}{\partial M}(\delta e)_1} = \frac{\mu \gamma \frac{1}{2} \frac{1}{2}}{c a^{2+1/2}} \sum_q \sum_s s a_{s,q} B_{n,q} B_{s,s} \sin \beta (\overline{E'_n \delta x p})$$

To find  $\overline{E'_n \delta x p}$  simply replace  $\sin \alpha$  in Equation (11) of the inclination equation by  $-\cos \alpha$ .

$$2. \quad \overline{\frac{\partial R}{\partial M}(\delta e)_2} = \frac{-\mu \gamma \frac{1}{2} \frac{1}{2}}{c a^{2+1/2}} \sum_q \sum_s \sum_{k=0}^{\infty} \epsilon_k s a_{s,q,k} B_{n,q} B_{s,s} (\overline{E'_n c e})$$

- To find  $\overline{E'_n c e}$  replace  $\sin(\beta - \alpha)$  in Equation (12) of the inclination equation by  $-\sin(\beta - \alpha)$ .

The Equation For the Longitude of the Node

The expansion for the node corresponding to Equation(3) is given by

$$\left[ \frac{d\Omega}{dt} \right] = \left[ \frac{1}{f} \left\{ a \frac{\partial^2 R}{\partial i \partial a} \frac{\delta a}{a} + \frac{\partial^2 R}{\partial a \partial e} \delta e + \frac{\partial^2 R}{\partial i^2} \delta i + \frac{\partial^2 R}{\partial a \partial \omega} \delta \omega \right. \right. \\ \left. \left. + \frac{\partial R}{\partial i \partial H} \delta H \right\} + \frac{\partial R}{\partial i} \left\{ a \frac{\partial f}{\partial a} \frac{\delta a}{a} + \frac{\partial f}{\partial e} \delta e + \frac{\partial f}{\partial i} \delta i \right\} \right]$$

$$f = \frac{1}{\mu^2 a^4 \eta \sin i}$$

$$a \frac{\partial f}{\partial a} = \frac{-1}{2 \mu^2 a^4 \eta \sin i}$$

$$\frac{\partial f}{\partial e} = \frac{e}{\mu^2 a^4 \eta^3 \sin i}$$

$$\frac{\partial f}{\partial i} = \frac{-\cos i}{\mu^2 a^4 \eta \sin^3 i}$$

When the products are formed it will be seen below that each of the long period factors of  $\left[ \frac{d\Omega}{dt} \right]$  are the conjugates of the corresponding long period factors of  $\left[ \frac{di}{dt} \right]$ . Consequently it is only necessary to replace each  $\alpha$  of the Equations 1 through 6, and 15 through 20, of  $\left[ \frac{di}{dt} \right]$  by  $\alpha + \frac{\pi}{2}$ , and each  $\alpha$  in equations 7 through 14, of  $\left[ \frac{d\Omega}{dt} \right]$  by  $\alpha - \frac{\pi}{2}$ , to obtain the corresponding equations for  $\left[ \frac{d\Omega}{dt} \right]$ .

$$1. \quad a \frac{\partial^2 R}{\partial a \partial a} \frac{\delta a}{a} = \frac{-2\mu(n+1) L_2 \sqrt{e} \cos i}{a^{2n+11}} \sum_q \sum_s B'_{n,q} B_{s,0} (\overline{A_n A_e} - \overline{A_n} \overline{A_e})$$

$$2. \quad \frac{\partial^2 R}{\partial i \partial e} (\delta e) = \frac{\mu \eta^2 L_2 \sqrt{e} \cos i}{e a^{n+11}} \sum_q \sum_s B'_{n,q} B_{s,0} (\overline{F_n A_e} - \overline{F_n} \overline{A_e})$$

$$3. \frac{\partial^2 R}{\partial i \partial e} (\delta e)_2 = \frac{\mu \eta J_1 J_2 \cos i}{e a^{n+l+1}} \sum_q \sum_s s a_{l,s} B'_{n,q} B_{l,s} \sin \beta (\overline{F_n} \delta_{kr})$$

$$4. \frac{\partial^2 R}{\partial i \partial e} (\delta e)_3 = \frac{-\mu \eta J_1 J_2 \cos i}{e a^{n+l+1}} \sum_q \sum_s \sum_{k=0}^{l-1} e_k s a_{l-1,k} B'_{n,q} B_{l,s} (\overline{F_n} \bar{C}_e - \overline{F_n} \bar{C}_e)$$

$$5. \frac{\partial^2 R}{\partial e^2} (\delta e)_1 = \frac{-\mu J_1 J_2 \cos i}{\eta a^{n+l+1} \sin i} \sum_q \sum_s s a_{l,s} (-\sin i B'_{n,q} + \cos i B'_{n,q}) B_{l,s} \sin \beta (\overline{A_n} \delta_{kr})$$

$$6. \frac{\partial^2 R}{\partial e^2} (\delta e)_2 = \frac{-\mu J_1 J_2 \cos i}{\eta a^{n+l+1} \sin i} \sum_q \sum_s \sum_{k=0}^{l-1} e_k s a_{l-1,k} (-\sin i B'_{n,q} + \cos i B'_{n,q}) B_{l,s} (\overline{A_n} \bar{C}_e - \overline{A_n} \bar{C}_e)$$

$$7. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_1 = \frac{-\mu J_1 J_2 \cos i}{\eta^2 e a^{n+l+1}} \sum_q \sum_s q B'_{n,q} W \cos \beta (\overline{A_n} \delta_{kr})$$

$$8. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_2 = \frac{-\mu J_1 J_2 \cos i}{\eta^2 e a^{n+l+1}} \sum_q \sum_s \sum_{k=0}^{l-1} e_k q B'_{n,q} W (\overline{A_n} \bar{C}'_e - \overline{A_n} \bar{C}'_e)$$

$$9. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_3 = \frac{-\mu J_1 J_2 \cos i}{2 e a^{n+l+1}} \sum_q \sum_s q B'_{n,q} B_{l,s} (\overline{A_n} \bar{D}'_{le} - \overline{A_n} \bar{D}'_{le})$$

$$10. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_4 = \frac{-\mu J_1 J_2 \cos i}{2 e \eta^2 a^{n+l+1}} \sum_q \sum_s q B'_{n,q} B_{l,s} (\overline{A_n} \bar{D}'_{2e} - \overline{A_n} \bar{D}'_{2e})$$

$$11. \frac{\partial^2 R}{\partial i \partial M} (\delta M)_1 = \frac{-\mu J_1 J_2 \cos i}{e \eta^{2l+3} a^{n+l+1}} \sum_q \sum_s l'_{l-1,s} B'_{n,q} B_{l,s} \cos \beta (\overline{E_n} \delta_{kr})$$

$$12. \frac{\partial^2 R}{\partial i \partial M} (\delta M)_2 = \frac{-\mu J_1 J_2 \cos i}{e \eta^{2l+3} a^{n+l+1}} \sum_q \sum_s \sum_{k=0}^{l-1} e_k l'_{l-1,k} B'_{n,q} B_{l,s} (\overline{E_n} \bar{C}'_e)$$



$$13. \frac{\partial^2 R}{\partial i \partial M} (\delta M)_3 = \frac{-\mu \eta J_n J_e \cos i}{e a^{n+1/2}} \sum_q \sum_s B'_{n,q} B_{e,s} (\overline{E'_n D'_{e,s}})$$

$$14. \frac{\partial^2 R}{\partial u \partial M} (\delta M)_4 = \frac{-\mu \eta J_n J_e \cos i}{e \eta a^{n+1/2}} \sum_q \sum_s B'_{n,q} B_{e,s} (\overline{E'_n D'_{e,s}})$$

$$15. \frac{\partial^2 R}{\partial e \partial a} = \frac{2\mu J_n J_e \cos i}{a^{n+1/2}} \sum_q \sum_s B'_{n,q} B_{e,s} (\overline{A'_n A_e} - \overline{A'_n \tilde{A}_e})$$

$$16. \frac{\partial^2 R}{\partial e} (\delta e)_1 = \frac{\mu \eta J_n J_e \cos i}{e a^{n+1/2}} \sum_q \sum_s B'_{n,q} B_{e,s} (\overline{A'_n A_e} - \overline{A'_n \tilde{A}_e})$$

$$17. \frac{\partial^2 R}{\partial e} (\delta e)_2 = \frac{\mu \eta J_n J_e \cos i}{e a^{n+1/2}} \sum_q \sum_s s a_{e,1,s} B'_{n,q} B_{e,s} \sin \beta (\overline{A'_n B_{xp}})$$

$$18. \frac{\partial^2 R}{\partial i} (\delta e)_3 = \frac{-\mu \eta J_n J_e \cos i}{e a^{n+1/2}} \sum_q \sum_s \sum_{k=0}^{e-1} e_k s a_{e,1,k} B'_{n,q} B_{e,s} (\overline{A'_n C_e} - \overline{A'_n \tilde{C}_e})$$

$$19. \frac{\partial^2 R}{\partial e} (\delta i)_1 = \frac{-\mu J_n J_e \cos i}{\eta a^{n+1/2} \sin i} \sum_q \sum_s s a_{e,1,s} B'_{n,q} B_{e,s} (\overline{A'_n B_{xp}})$$

$$20. \frac{\partial^2 R}{\partial i} (\delta i)_2 = \frac{\mu J_n J_e \cos i}{\eta a^{n+1/2} \sin i} \sum_q \sum_s \sum_{k=0}^{e-1} e_k s a_{e,1,k} B'_{n,q} B_{e,s} (\overline{A'_n C_e} - \overline{A'_n \tilde{C}_e})$$

The long period terms for the equation of the node are the sum of the above 20 equations.

## The Argument of Perigee Equation

The expansion for the argument of perigee equation corresponding to Equation (3) is given by

$$\begin{aligned} \left[ \frac{d\omega}{dt} \right] = & \left[ f_{v1} \left\{ \left[ \frac{\partial R}{\partial i} \right] \right\} + \frac{\partial R}{\partial i} \left\{ a \frac{\partial f_{v1}}{\partial a} \frac{\delta a}{a} + \frac{\partial f_{v1}}{\partial e} \delta e + \frac{\partial f_{v1}}{\partial i} \delta i \right\} \right. \\ & + f_{v2} \left\{ a \frac{\partial^2 R}{\partial e \partial a} \frac{\delta a}{a} + \frac{\partial^2 R}{\partial e^2} \delta e + \frac{\partial^2 R}{\partial e \partial i} \delta i + \frac{\partial^2 R}{\partial e \partial \omega} \delta \omega + \frac{\partial^2 R}{\partial e \partial M} \delta M \right\} \\ & \left. + \frac{\partial R}{\partial e} \left\{ a \frac{\partial f_{v2}}{\partial a} \frac{\delta a}{a} + \frac{\partial f_{v2}}{\partial e} \delta e \right\} \right] \end{aligned}$$

$$f_{v1} = - \frac{\cos i}{\mu^2 a^2 \eta \sin i}$$

$$f_{v2} = \frac{\eta}{\mu^2 a^2 e}$$

$$a \frac{\partial f_{v1}}{\partial a} = \frac{\cos i}{2 \mu^2 a^2 \eta \sin i}$$

$$a \frac{\partial f_{v2}}{\partial a} = - \frac{\eta}{2 \mu^2 a^2 e}$$

$$\frac{\partial f_{v1}}{\partial e} = - \frac{e \cos i}{\mu^2 a^2 \eta^3 \sin i}$$

$$\frac{\partial f_{v2}}{\partial e} = \frac{-1}{\mu^2 a^2 e^2 \eta}$$

$$\frac{\partial f_{v1}}{\partial i} = \frac{1}{\mu^2 a^2 \eta \sin^2 i}$$

The long period terms for the expansion  $\left\{ \left[ \frac{\partial R}{\partial i} \right] \right\}$ , where

$$\begin{aligned} \left\{ \left[ \frac{\partial R}{\partial i} \right] \right\} = & a \frac{\partial^2 R}{\partial i \partial a} \frac{\delta a}{a} + \frac{\partial^2 R}{\partial i \partial e} \delta e + \frac{\partial^2 R}{\partial i^2} \delta i + \frac{\partial^2 R}{\partial i \partial \omega} \delta \omega \\ & + \frac{\partial^2 R}{\partial i \partial M} \delta M, \end{aligned}$$

as well as the long period terms of the functions  $\frac{\partial R}{\partial a} \frac{\delta a}{a}$ ,  $\frac{\partial R}{\partial e} \delta e$ ,

$\frac{\partial R}{\partial i} \delta i$  are given in the equation for the node. In addition since the long period terms  $(\overline{F_n'} A_e - \overline{F_n'} \overline{A_e})$ ,  $\overline{F_n'} B_{xp}$ ,  $\overline{F_n'} C_e - \overline{F_n'} \overline{C_e}$ , are given in the inclination equation, the corresponding terms of the perigee equation may be derived by replacing  $\kappa$  in the above terms of the inclination equation by  $\kappa + \frac{\pi}{2}$ .



$$1. \frac{\partial^2 R}{\partial \epsilon \partial \alpha} \frac{\partial \alpha}{\partial \epsilon} = \frac{-2\mu(n+1) \frac{1}{2} J_0}{a^{n+1}} \sum_q \sum_s B_{n,q} B_{s,0} (\overline{F_n A_s} - \overline{F_n} \overline{A_s})$$

$$2. \frac{\partial^2 R(\delta \epsilon)_1}{\partial \epsilon^2} = \frac{\mu \eta^2 \frac{1}{2} J_0}{\epsilon a^{n+1}} \sum_q \sum_s B_{n,q} B_{s,0} (\overline{H_n A_s} - \overline{H_n} \overline{A_s})$$

$$\overline{H_n A_s} = \sum_{d=1}^9 \frac{c(d)}{2} \{ \cos(\alpha-\beta) a_{\nu(d), \delta(d)} + \cos(\alpha+\beta) a_{\nu(d), \rho(d)} \}$$

The parameters for  $\overline{H_n A_s}$  are defined in Table 7.

TABLE 7 ( Parameters for  $\overline{H_n A_s}$  )

d	c(d)	v(d)	δ(d)	ρ(d)
1	$\frac{(n+q+1)(n+q+3)}{4}$	$n+l+2$	$q-s+2$	$q+s+2$
2	$\frac{2nq+2q^2+4q+n+1}{4\eta^2}$	$n+l+1$	$q-s+2$	$q+s+2$
3	$\frac{q(q+1)}{4\eta^2}$	$n+l$	$q-s+2$	$q+s+2$
4	$\frac{(n-q+1)(n-q+3)}{4}$	$n+l+2$	$q-s-2$	$q+s-2$
5	$\frac{n+1+2q^2+2nq+4q}{4\eta^2}$	$n+l+1$	$q-s-2$	$q+s-2$
6	$\frac{q(q-1)}{4\eta^2}$	$n+l$	$q-s-2$	$q+s-2$
7	$\frac{(n+q+1)(n-q+1)}{4}$	$n+l+2$	$q-s$	$q+s$
8	$\frac{(2q^2+n+1)}{2\eta^2}$	$n+l+1$	$q-s$	$q+s$
9	$-\frac{q^2}{2\eta^2}$	$n+l$	$q-s$	$q+s$

$$\overline{H_n A_s} = \left\{ \frac{h''_{n+1,q}}{\eta^{2n+1}} - \frac{(2n+1)}{\eta} \left[ \frac{a_{n+1,q}}{\eta} \left( \frac{\epsilon}{\eta} - 2n+1 \right) + 2 h_{n+1,q} \right] \right\} a_{\nu(s), \delta(s)} \cos \alpha \cos \beta$$

$$3. \frac{\partial^2 K}{\partial c^2} (\delta e)_n = \frac{\mu \eta \frac{1}{2} l_0}{c a^{n+1}} \sum_{\substack{q \\ \neq}} \sum_{\substack{s \\ \neq}} s a_{n,q} b_{n,q} \sin \beta (\overline{H} \overline{B}_{kp})$$

$$\overline{H}_n \overline{B}_{kp} = \sum_{d=1}^9 c(d) \left\{ \sum_{t=1}^{t_1} a_{n(d),q(d)} - \sum_{t=2}^{t_2} a_{n(d),f(d)} \right\} \sin \alpha$$

The parameters for  $\overline{H}_n \overline{B}_{kp}$  are defined in Table 8.

TABLE 8 ( Parameters for  $\overline{H}_n \overline{B}_{kp}$ ,  $c(d)$ ,  $q(d)$ , same as Table 7)

$d$	$\delta(d)$	$t_1$	$t(d)$	$t_2$
1	$q-t+2$	$n+l+q+4$	$q+t+2$	$n+l-q$
2	$q-t+2$	$n+l+q+3$	$q+t+2$	$n+l-q-1$
3	$q-t+2$	$n+l+q+2$	$q+t+2$	$n+l-q-2$
4	$q-t-2$	$n+l+q$	$q+t-2$	$n+l-q+4$
5	$q-t-2$	$n+l+q-1$	$q+t-2$	$n+l-q+3$
6	$q-t-2$	$n+l+q-2$	$q+t-2$	$n+l-q+2$
7	$q-t$	$n+l+q+2$	$q+t$	$n+l-q+2$
8	$q-t$	$n+l+q+1$	$q+t$	$n+l-q+1$
9	$q-t$	$n+l+q$	$q+t$	$n+l-q$

$$4. \frac{\partial^2 K}{\partial c^2} (\delta e)_s = \frac{-\mu \eta \frac{1}{2} l_0}{c a^{n+1}} \sum_{\substack{q \\ \neq}} \sum_{\substack{s \\ \neq}} \sum_{k=0}^{k_1} \epsilon_k s a_{n,k} b_{n,q} B_{s,q} (\overline{H}_n \overline{C}_e - \overline{H}_n \overline{C}_e)$$

$$\overline{H}_n \overline{C}_e = \sum_{d=1}^9 \frac{c(d)}{2} \left\{ \frac{\cos(n-\delta)}{s+k} a_{n(d),q(d)} + \frac{\cos(n+\delta)}{s+k} a_{n(d),f(d)} + \frac{\cos(n-\delta)}{s-k} a_{n(d),\sigma(d)} + \frac{\cos(n+\delta)}{s-k} a_{n(d),\tau(d)} \right\}$$

The parameters for  $\overline{H}_n \overline{C}_e$  are defined in Table 9.

TABLE 9 ( Parameters for  $\overline{H}_n \overline{C}_e$ ,  $c(d)$  same as Table 7)

$d$	$v(d)$	$\delta(d)$	$t(d)$	$\sigma(d)$	$\tau(d)$
1	$n+1$	$q-s-k+2$	$q+s+k+2$	$q-s+k+2$	$q+s-k+2$
2	$n$	$q-s-k+2$	$q+s+k+2$	$q-s+k+2$	$q+s-k+2$
3	$n-1$	$q-s-k+2$	$q+s+k+2$	$q-s+k+2$	$q+s-k+2$
4	$n+1$	$q-s-k-2$	$q+s+k-2$	$q-s+k-2$	$q+s-k-2$
5	$n$	$q-s-k-2$	$q+s+k-2$	$q-s+k-2$	$q+s-k-2$
6	$n-1$	$q-s-k-2$	$q+s+k-2$	$q-s+k-2$	$q+s-k-2$
7	$n+1$	$q-s-k$	$q+s+k$	$q-s+k$	$q+s-k$
8	$n$	$q-s-k$	$q+s+k$	$q-s+k$	$q+s-k$
9	$n-1$	$q-s-k$	$q+s+k$	$q-s+k$	$q+s-k$

$$\overline{H}_n \overline{C}_e = a_{n+2} (r_{n+2} - r_{0-2}) \left\{ \frac{a_{n+2}}{r_{n+2}} - \frac{(2n+1)}{n} \left[ \frac{a_{n+2}}{n} \left( \frac{c}{n} - 2n+1 \right) + 2 \frac{b_{n+2}}{n} \right] \right\} \cos \alpha \cos \beta$$

$$5. \frac{\partial^2 R}{\partial e \partial e} (\delta e)_1 = \frac{-\mu J_n J_e \cos i}{\eta^2 a^{2+2H} \sin i} \sum_q \sum_s s a_{2s,q} B'_{n,q} B_{e,s} \sin \beta (\overline{F_n} \overline{B_{kp}})$$

$$6. \frac{\partial^2 R}{\partial e \partial e} (\delta e)_2 = \frac{\mu J_n J_e \cos i}{\eta^2 a^{2+2H} \sin i} \sum_q \sum_s \sum_{k=0}^{l-1} \epsilon_k s a_{2s,q} B'_{n,q} B_{e,s} (\overline{F_n} \overline{C_e} - \overline{F_n} \overline{C_e})$$

$$7. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_1 = \frac{-\mu J_n J_e}{\eta^2 a^{2+2H}} \sum_q \sum_s q B_{n,q} W \cos \beta (\overline{F_n} \overline{B_{kp}})$$

$$8. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_2 = \frac{-\mu J_n J_e}{\eta^2 a^{2+2H}} \sum_q \sum_s \sum_{k=0}^{l-1} q \epsilon_k B_{n,q} W (\overline{F_n} \overline{C'_e} - \overline{F_n} \overline{C'_e})$$

$$9. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_3 = \frac{-\mu J_n J_e}{2e a^{2+2H}} \sum_q \sum_s q B_{n,q} B_{e,s} (\overline{F_n} \overline{D'_{1e}} - \overline{F_n} \overline{D'_{1e}})$$

$$\overline{F_n} \overline{D'_{1e}} = \sum_{d=1}^4 \frac{c(d)}{2} \{ \cos(\nu - \beta) [a_{\nu(d), \phi(d)} - a_{\nu(d), \phi(d)}] + \cos(\nu + \beta) [a_{\nu(d), \sigma(d)} - a_{\nu(d), \sigma(d)}] \}$$

The parameters for  $\overline{F_n} \overline{D'_{1e}}$  are defined in Table 10.

TABLE 10 (Parameters for  $\overline{F_n} \overline{D'_{1e}}$ )

d	c(d)	$\nu(d)$	$\phi(d)$	$\sigma(d)$	$\sigma(d)$	$\sigma(d)$
1	$\frac{n+1+q}{2}$	$n+l$	$q-s$	$q-s+2$	$q+s+2$	$q+s$
2	$\frac{n+1+q}{2}$	$n+l$	$q-s-2$	$q-s$	$q+s$	$q+s-2$
3	$\frac{q}{2\eta^2}$	$n+l-1$	$q-s$	$q-s+2$	$q+s+2$	$q+s$
4	$-\frac{q}{2\eta^2}$	$n+l-1$	$q-s-2$	$q-s$	$q+s$	$q+s-2$

$$\overline{F_n} \overline{D'_{1e}} = a_{1,2,1+1} [t_{n,q} - \frac{(2n-1)}{\eta} a_{n-1,q}] \sin \alpha \sin \beta$$

$$10. \frac{\partial^2 R}{\partial e \partial \omega} (\delta \omega)_4 = \frac{-\mu J_n J_e}{2e \eta^2 a^{2+2H}} \sum_q \sum_s q B_{n,q} B_{e,s} (\overline{F_n} \overline{D'_{1e}} - \overline{F_n} \overline{D'_{1e}})$$

$$\overline{F_n} \overline{D'_{1e}} = \sum_{d=1}^4 \frac{c(d)}{2} \{ \cos(\alpha - \beta) [a_{\nu(d), \phi(d)} - a_{\nu(d), \phi(d)}] + \cos(\alpha + \beta) [a_{\nu(d), \sigma(d)} - a_{\nu(d), \sigma(d)}] \}$$

$$\overline{F}_n' = \sin \alpha \left[ \overline{G}_{n+1, q} - \frac{(2n+1)}{\eta} \overline{a}_{n+1, q} \right]$$

$$\overline{D}_{2c}' = \sin \beta \left[ \overline{a}_{2,3,3,1} - \overline{a}_{2,3,4,1} \right] \quad \text{for } l \geq 3$$

$$\overline{D}_{2c}' = \frac{\sin \beta}{\eta^2} \left\{ \overline{\cos}(\overline{a}_{1,1}) - \overline{\cos}(\overline{a}_{1,2}) + \overline{e} \left( \overline{\cos}(\overline{a}_{1,2}) - \overline{\cos}(\overline{a}_{1,3}) \right) \right\} \quad \text{for } l=2$$

The definition of  $\overline{\cos}$  is given in appendix 1. The parameters  $v(d)$  for  $\overline{F}_n' \overline{D}_{2c}'$  are defined in Table 11. The remaining parameters of  $\overline{F}_n' \overline{D}_{2c}'$  are the same as those defined in Table 10.

TABLE 11 (Parameters  $v(d)$  for  $\overline{F}_n' \overline{D}_{2c}'$ )

$d$	$v(d)$
1	$n+l-1$
2	$n+l-1$
3	$n+l-2$
4	$n+l-2$

$$11. \frac{\partial^2 R}{\partial \epsilon \partial M} (\delta H)_1 = \frac{-\mu \ln \epsilon}{\eta^{2.5} \epsilon a^{n+l-1}} \sum_q \sum_t \overline{G}_{n+1, q}' B_{n, q} B_{t, c} \cos \beta (\overline{G}_n' B_{kp})$$

$$\overline{G}_n' B_{kp} = \sum_{d=1}^{10} c(d) \left\{ \sum_{t=1}^{t_1} \overline{r}_t \overline{a}_{n+l, q}(d) - \sum_{t=1}^{t_2} \overline{a}_{n+l, q}(d) \right\} \cos \alpha$$

The parameters of  $\overline{G}_n' B_{kp}$  are given in Table 12.

TABLE 12 (Parameters of  $\overline{G}_n' B_{kp}$ )

$d$	$c(d)$	$v(d)$	$\phi(d)$	$t_1$	$t_2$	$x_2$
1	$\frac{e(n+q)(n+2)}{4\eta}$	$n+1$	$q+t+2$	$n-q-1$	$q-t+2$	$n+q+3$
2	$\frac{e(n+1)}{4\eta^2}$	$n$	$q+t+2$	$n-q-2$	$q-t+2$	$n+q+2$
3	$-\frac{\eta(n+q)(q+1)}{2}$	$n+2$	$q+t+1$	$n-q+1$	$q-t+1$	$n+q+3$
4	$-\frac{e(q+1)}{2\eta}$	$n+1$	$q+t+1$	$n-q$	$q-t+1$	$n+q+2$
5	$\frac{e q(n+2)}{2\eta}$	$n+1$	$q+t$	$n-q+1$	$q-t$	$n+q+1$
6	$\frac{e q(n+1)}{2\eta^2}$	$n$	$q+t$	$n-q$	$q-t$	$n+q$
7	$-\frac{\eta(n+1-q)(q-2)}{2}$	$n+2$	$q+t-1$	$n-q+3$	$q-t-1$	$n+q+1$
8	$\frac{q(q+1)}{2\eta}$	$n+1$	$q+t-1$	$n-q+2$	$q-t-1$	$n+q$
9	$\frac{e(n+1-q)(n+2)}{4\eta}$	$n+1$	$q+t-2$	$n-q+3$	$q-t-2$	$n+q-1$
10	$-\frac{e(n+1)q}{4\eta^2}$	$n$	$q+t-2$	$n-q+2$	$q-t-2$	$n+q-2$



$$12. \frac{\partial^2 R}{\partial c \partial H} (\bar{S}H)_c = \frac{-u \gamma \ln k}{c \eta^{l+3} a^{n+l+1}} \sum_{\frac{1}{2}} \sum_{\frac{1}{2}} \sum_{k=0}^{l+1} \epsilon_k \epsilon'_{l+1-k} S_{n,\frac{1}{2}} S_{l,\frac{1}{2}} (\bar{G}'_n \bar{C}'_c)$$

$$\bar{G}'_n \bar{C}'_c = \sum_{d=1}^{10} \frac{c(d)}{2} \left\{ \frac{\cos(\alpha-\beta)}{s+k} [a_{v(d),\phi(d)} - a_{v(d),\psi(d)}] + \frac{\cos(\beta-\alpha)}{s+k} [a_{v(d),\psi(d)} + \frac{\cos(\alpha-\beta)}{s-k} a_{v(d),\psi(d)}] - \frac{\cos(\alpha-\beta)}{s-k} a_{v(d),\psi(d)} \right\}$$

The parameters of  $\bar{G}'_n \bar{C}'_c$  are given in Table 13.

TABLE 13 (Parameters of  $\bar{G}'_n \bar{C}'_c$ ,  $c(d)$ ,  $v(d)$ , same as Table 11)

d	$\phi(d)$	$\psi(d)$	$v(d)$	$\tau(d)$
1	$q-s-k+2$	$q+s+k+2$	$q-s+k+2$	$q+s-k+2$
2	$q-s-k+2$	$q+s+k+2$	$q-s+k+2$	$q+s-k+2$
3	$q-s-k+1$	$q+s+k+1$	$q-s+k+1$	$q+s-k+1$
4	$q-s-k+1$	$q+s+k+1$	$q-s+k+1$	$q+s-k+1$
5	$q-s-k$	$q+s+k$	$q-s+k$	$q+s-k$
6	$q-s-k$	$q+s+k$	$q-s+k$	$q+s-k$
7	$q-s-k-1$	$q+s+k-1$	$q-s+k-1$	$q+s-k-1$
8	$q-s-k-1$	$q+s+k-1$	$q-s+k-1$	$q+s-k-1$
9	$q-s-k-2$	$q+s+k-2$	$q-s+k-2$	$q+s-k-2$
10	$q-s-k-2$	$q+s+k-2$	$q-s+k-2$	$q+s-k-2$

$$13. \frac{\partial^2 R}{\partial c \partial H} (\bar{S}H)_c = \frac{-u \gamma \ln k}{c a^{n+l+1}} \sum_{\frac{1}{2}} \sum_{\frac{1}{2}} B_{n,\frac{1}{2}} B_{l,\frac{1}{2}} (\bar{G}'_n \bar{D}'_c)$$

$$\bar{G}'_n \bar{D}'_c = \sum_{d=1}^{10} \frac{c(d)}{4} \left\{ \cos(\alpha-\beta) [a_{v(d),\phi(d)} - a_{v(d),\psi(d)}] - \cos(\beta-\alpha) [a_{v(d),\psi(d)} - a_{v(d),\tau(d)}] \right\}$$

The parameters of  $\bar{G}'_n \bar{D}'_c$  are given in Table 14.

TABLE 14 (Parameters of  $\bar{G}'_n \bar{D}'_c$ ,  $c(d)$ , same as Table 11)

d	$v(d)$	$\phi(d)$	$\psi(d)$	$\tau(d)$	$\tau(d)$
1	$n+l+1$	$q-s+1$	$q-s+3$	$q+s+3$	$q+s+1$
2	$n+l$	$q-s+1$	$q-s+3$	$q+s+3$	$q+s+1$
3	$n+l+2$	$q-s$	$q-s+2$	$q+s+2$	$q+s$
4	$n+l+1$	$q-s$	$q-s+2$	$q+s+2$	$q+s$
5	$n+l+1$	$q-s-1$	$q-s+1$	$q+s+1$	$q+s-1$
6	$n+l$	$q-s-1$	$q-s+1$	$q+s+1$	$q+s-1$
7	$n+l+2$	$q-s-2$	$q-s$	$q+s$	$q+s-2$
8	$n+l+1$	$q-s-2$	$q-s$	$q+s$	$q+s-2$
9	$n+l+1$	$q-s-3$	$q-s-1$	$q+s-1$	$q+s-3$
10	$n+l+2$	$q-s-3$	$q-s-1$	$q+s-1$	$q+s-3$

$$14. \frac{\partial R}{\partial e} (\overline{G_{1e}}) = \frac{-\mu \gamma \gamma_e}{e \gamma a^{n+l+1}} \sum_q \sum_s B_{n,q} B_{e,s} (\overline{G_n} \overline{D'_{1e}})$$

$$\overline{G_n} \overline{D'_{1e}} = \sum_{d=1}^4 \frac{c(d)}{4} \{ \cos(\omega+\beta) [a_{nd}, q(d) \cdot a_{nd}, r(d)] - \cos(\omega+\beta) [a_{nd}, q(d) \cdot a_{nd}, r(d)] \}$$

The parameters  $\frac{c(d)}{4}$  for  $\overline{F_n} \overline{D'_{1e}}$  are defined in Table 15. The remaining parameters of  $\overline{F_n} \overline{D'_{1e}}$  are the same as those defined in Table 14.

TABLE 15 (Parameters  $v(d)$  for  $\overline{F_n} \overline{D'_{1e}}$ )

d	v(d)
1	n+l
2	n+l-1
3	n+l+1
4	n+l
5	n+l
6	n+l-1
7	n+l+1
8	n+l
9	n+l
10	n+l+1

$$15. \frac{\partial R}{\partial e} \frac{\overline{A_e}}{a} = \frac{2\mu \gamma \gamma_e}{e a^{n+l+1}} \sum_q \sum_s B_{n,q} B_{e,s} (\overline{F_n} \overline{A_e} - \overline{F_n} \overline{A_e})$$

$$16. \frac{\partial R}{\partial e} (\overline{D_e})_1 = \frac{\mu \gamma \gamma_e}{e a^{n+l+1}} \sum_q \sum_s B_{n,q} B_{e,s} (\overline{F_n} \overline{A_e} - \overline{F_n} \overline{A_e})$$

$$17. \frac{\partial R}{\partial e} (\overline{D_e})_2 = \frac{\mu \gamma \gamma_e}{e a^{n+l+1}} \sum_q \sum_s s a_{e,s} B_{n,q} B_{e,s} \sin \beta (\overline{F_n} \overline{B_{ep}})$$

$$18. \frac{\partial R}{\partial e} (\overline{D_e})_3 = \frac{-\mu \gamma \gamma_e}{e a^{n+l+1}} \sum_q \sum_s \sum_{k=0}^{e-1} c_k s a_{e,k} B_{n,q} B_{e,s} (\overline{F_n} \overline{C_e} - \overline{F_n} \overline{C_e})$$

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## The Equation For the Mean Anomaly

The expansion for the mean anomaly corresponding to Equation (3) is given by

$$\begin{aligned} \left[ \frac{dM}{dt} \right] = & \left[ \frac{15}{8} \mu^2 a^{-2} \left( \frac{\delta a}{a} \right)_n \left( \frac{\delta a}{a} \right)_i - f_{61} \left\{ \left[ \frac{\partial R}{\partial e} \right] \right\} - \frac{\partial R}{\partial e} \left\{ a \frac{\partial f_{61}}{\partial a} \frac{\delta a}{a} + \frac{\partial f_{61}}{\partial e} \delta e \right\} \right. \\ & - \frac{f_{62}}{a} \left\{ a^2 \frac{\partial^2 R}{\partial a^2} \frac{\delta a}{a} + a \frac{\partial^2 R}{\partial a \partial e} \delta e + a \frac{\partial^2 R}{\partial a \partial i} \delta i + a \frac{\partial^2 R}{\partial a \partial \omega} \delta \omega + a \frac{\partial^2 R}{\partial a \partial M} \delta M \right\} \\ & \left. - a \frac{\partial R}{\partial a} \left\{ \frac{\partial f_{61}}{\partial a} \frac{\delta a}{a} \right\} \right]. \end{aligned}$$

The first product  $\left( \frac{\delta a}{a} \right)_n \left( \frac{\delta a}{a} \right)_i$  results from the expansion  $a^{-2}$  in powers of  $\frac{\delta a}{a}$ . The quantities

$$\left\{ \left[ \frac{\partial R}{\partial e} \right] \right\} = a \frac{\partial^2 R}{\partial a \partial e} \frac{\delta a}{a} + \frac{\partial^2 R}{\partial e^2} \delta e + \frac{\partial^2 R}{\partial e \partial i} \delta i + \frac{\partial^2 R}{\partial e \partial \omega} \delta \omega + \frac{\partial^2 R}{\partial e \partial M} \delta M$$

$a \frac{\partial^2 R}{\partial a} \frac{\delta a}{a}$ ,  $\frac{\partial R}{\partial e} \delta e$ , have been already defined in the argument of perigee equation. The remaining products of  $\left[ \frac{dM}{dt} \right]$  are given below. These products are the conjugates of the products given in  $\left[ \frac{di}{dt} \right]$ . Thus the products in equations 1 through 6, and Equation (15), of  $\left[ \frac{dM}{dt} \right]$  are found by replacing  $\alpha$  of the corresponding equations of  $\left[ \frac{di}{dt} \right]$  by  $\alpha + \frac{\pi}{2}$ . Similarly Equations 7 through 14 of  $\left[ \frac{dM}{dt} \right]$  are found by replacing  $\alpha$  of the corresponding equations in  $\left[ \frac{di}{dt} \right]$  by  $\alpha - \frac{\pi}{2}$ .

$$f_{61} = \frac{\eta^2}{u^2 a^2 e}$$

$$\frac{f_{62}}{a} = \frac{2}{u^2 a^2}$$

$$a \frac{\partial f_{61}}{\partial a} = \frac{-\eta^2}{u^2 a^2 e}$$

$$\frac{\partial f_{62}}{\partial a} = \frac{1}{u^2 a^2}$$

$$\frac{\partial f_{61}}{\partial e} = \frac{-(1+e^2)}{u^2 a^2 e^2}$$

$$1. \frac{a^2 \delta^2 R}{\delta a^2} \frac{\delta a}{a} = \frac{2\mu(n+1)(n+2)J_n J_e}{a^{n+2e+1}} \sum_q \sum_s B_{n,q} B_{e,s} (\overline{A_n A_e} - \overline{A_n} \overline{A_e})$$

$$2. \frac{a \delta^2 R}{\delta a \delta e} (\delta e)_1 = -\frac{\mu(n+1)\eta^2 J_n J_e}{e a^{n+2e+1}} \sum_q \sum_s B_{n,q} B_{e,s} (\overline{F_n A_e} - \overline{F_n} \overline{A_e})$$

$$3. \frac{a \delta^2 R}{\delta a \delta e} (\delta e)_2 = -\frac{\mu(n+1)\eta J_n J_e}{e a^{n+2e+1}} \sum_q \sum_s s a_{e-1,s} B_{n,q} B_{e,s} \sin \beta (\overline{F_n B_{kp}})$$

$$4. \frac{a \delta^2 R}{\delta a \delta e} (\delta e)_3 = \frac{\mu(n+1)\eta J_n J_e}{e a^{n+2e+1}} \sum_q \sum_s \sum_{k=0}^{e-1} \epsilon_k s a_{e-1,k} B_{n,q} B_{e,s} (\overline{F_n C_e} - \overline{F_n} \overline{C_e})$$

$$5. \frac{a \delta^2 R}{\delta a \delta i} (\delta i)_1 = \frac{\mu(n+1)J_n J_e \cos^2 i}{\eta a^{n+2e+1} \sin i} \sum_q \sum_s s a_{e-1,s} B'_{n,q} B_{e,s} \sin \beta (\overline{A_n B_{kp}})$$

$$6. \frac{a \delta^2 R}{\delta a \delta i} (\delta i)_2 = -\frac{\mu(n+1)J_n J_e \cos^2 i}{\eta a^{n+2e+1} \sin i} \sum_q \sum_s \sum_{k=0}^{e-1} \epsilon_k s a_{e-1,k} B'_{n,q} B_{e,s} (\overline{A_n C_e} - \overline{A_n} \overline{C_e})$$

$$7. \frac{a \delta^2 R}{\delta a \delta \omega} (\delta \omega)_1 = \frac{\mu(n+1)J_n J_e}{\eta^2 a^{n+2e+1}} \sum_q \sum_s q B_{n,q} W \cos \beta (\overline{A_n B_{kp}})$$

$$8. \frac{a \delta^2 R}{\delta a \delta \omega} (\delta \omega)_2 = \frac{\mu(n+1)J_n J_e}{\eta^2 a^{n+2e+1}} \sum_q \sum_s \sum_{k=0}^{e-1} \epsilon_k q B_{n,q} W (\overline{A_n C_e} - \overline{A_n} \overline{C_e})$$

$$9. \frac{a \delta^2 R}{\delta a \delta \omega} (\delta \omega)_3 = \frac{\mu(n+1)J_n J_e}{2e a^{n+2e+1}} \sum_q \sum_s q B_{n,q} B_{e,s} (\overline{A_n D'_e} - \overline{A_n} \overline{D'_e})$$

$$10. \frac{a \delta^2 R}{\delta a \delta \omega} (\delta \omega)_4 = \frac{\mu(n+1)J_n J_e}{2e \eta^2 a^{n+2e+1}} \sum_q \sum_s q B_{n,q} B_{e,s} (\overline{A_n D'_e} - \overline{A_n} \overline{D'_e})$$

$$11. \frac{a \delta^2 R}{\delta a \delta H} (\delta H)_1 = \frac{\mu(n+1)J_n J_e}{e \eta^{2e+3} a^{n+2e+1}} \sum_q \sum_s t'_{e-1,s} B_{n,q} B_{e,s} \cos \beta (\overline{E_n B_{kp}})$$

$$12. \frac{a}{\partial a} \frac{\partial^2 R}{\partial M} (\delta H)_n = \frac{\mu(n+1) J_n J_e}{\eta^{2l+3} e^{a^{n+l+1}}} \sum_q \sum_s \sum_{k=0}^{l-1} \epsilon_k \delta'_{k,q} B_{n,q} B_{l,s} (\overline{E'_n C'_e})$$

$$13. \frac{a}{\partial a} \frac{\partial^2 R}{\partial M} (\delta H)_3 = \frac{\mu \eta(n+1) J_n J_e}{e^{a^{n+l+1}}} \sum_q \sum_s B_{n,q} B_{l,s} (\overline{E'_n D'_{le}})$$

$$14. \frac{a}{\partial a} \frac{\partial^2 R}{\partial M} (\delta H)_4 = \frac{\mu(n+1) J_n J_e}{e \eta^{2l+1}} \sum_q \sum_s B_{n,q} B_{l,s} (\overline{E'_n D'_{le}})$$

$$15. \frac{a}{\partial a} \frac{\partial^2 R}{\partial a} \frac{\partial}{\partial a} = -\frac{2\mu(n+1) J_n J_e}{a^{n+l+1}} \sum_q \sum_s B_{n,q} B_{l,s} (\overline{A_n A_e} - \overline{A_n} \overline{A_e})$$

$$16. \left( \frac{\partial a}{\partial a} \right)_n \left( \frac{\partial a}{\partial a} \right)_e = \frac{4 J_n J_e}{a^{n+l+1}} \sum_q \sum_s B_{n,q} B_{l,s} (\overline{A_n A_e} - \overline{A_n} \overline{A_e})$$

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## Appendix 1

## The Inclination Function

A convenient form of the inclination function is obtained from Reference 3.

$$B_{n,q} = \frac{1}{2^{n-1}} \cos[(q - ODQ) \frac{\pi}{2}] \epsilon_q \sum_{d=\frac{q-ODQ}{2}}^{\frac{n-ODQ}{2}} (-1)^{\frac{n-ODQ-2d}{2}} X$$

$$\left( d - \frac{q-ODQ}{2} \right) \frac{(n+2d+ODQ)! (\sin i)^{2d+ODQ}}{\left( \frac{n+2d+ODQ}{2} \right)! \left( \frac{n-2d-ODQ}{2} \right)! (2d+ODQ)! 2^{2d+ODQ}}$$

where  $ODQ = 0$  for  $q$  even or zero

$ODQ = 1$  for  $q$  odd

$$q = 2i + ODQ$$

$$i = 0, 1, 2, \dots, \frac{n-ODQ}{2}$$

$$B'_{n,q} = \frac{d B_{n,q}}{d(\sin i)}$$

$$B''_{n,q} = \frac{d^2 B_{n,q}}{d(\sin i)^2}$$

### The Averaging Function

The averaging function arises by first considering the expansion for

$$\left(\frac{a}{N}\right)^k = \frac{2}{\pi^{1/2}} \sum_{k=0}^{\infty} \epsilon_k a'_{k,k} \cos k f \quad \begin{matrix} \epsilon_k = \frac{1}{2} \text{ for } k=0, \\ = 1 \text{ for } k>0 \end{matrix} \quad (1)$$

$$a'_{k,k} = \sum_{d=0}^{k-k/2} \binom{2d+k}{d} \binom{k}{2d+k} \left(\frac{a}{2}\right)^{2d+k}$$

From the definition of the mean value, the function  $A$  is given by

$$\bar{A} = \frac{1}{2\pi} \int_0^{2\pi} \left(\frac{a}{N}\right)^{p''} \cos(\kappa + \theta \frac{f}{f}) dM$$

If we substitute Equation (1) and the expression for  $dM$  in terms of  $df$  we find

$$\left(\frac{a}{N}\right)^{p''} \cos(\kappa + \theta \frac{f}{f}) = a_{p-1,q}$$

where

$$a_{p,q} = \frac{a'_{p+1,q}}{\pi^{1/2}} \quad \text{for } q \leq p-1$$

$$a_{p-1,q} = 0 \quad q > p-1$$

$$a_{p-1,-q} = a_{p-1,q}$$

$$a_{0,q} = 0 \quad \text{for } q \neq 0$$

Expressions for the derivatives of the coefficients with respect, to  $e$  are also useful. We find

$$b'_{p-1,q} = \frac{d a_{p-1,q}}{d e} = \sum_{d=0}^{p+q-1} \frac{2d+q}{2} \binom{2d+q}{d} \binom{p-1}{2d+q} \left(\frac{e}{2}\right)^{2d+q-1}$$

$$b''_{p-1,q} = \frac{d^2 a_{p-1,q}}{d e^2} = \sum_{d=0}^{p+q-1} \frac{(2d+q)(2d+q-1)}{4} \binom{2d+q}{d} \binom{p-1}{2d+q} \left(\frac{e}{2}\right)^{2d+q-2}$$

$$b_{p-1,q} = \frac{b'_{p-1,q}}{\eta^{2p-1}}$$

$$\frac{d a_{p-1,q}}{d e} = b_{p-1,q} - \frac{(2p-1)}{\eta} a_{p-1,q}$$

$$\frac{d^2 a_{p-1,q}}{d e^2} = \frac{b''_{p-1,q}}{\eta^{2p-1}} - \frac{(2p-1)}{\eta} \left\{ \frac{a_{p-1,q}}{\eta} \left[ \frac{e}{\eta} - (2p-1) \right] + 2 b_{p-1,q} \right\}$$

The Equation of the Center

The expansion for the equation of the center (f-M) is given by Reference 1,

$$B_{KE} = f - M = -2 \sum_{p=1}^{\infty} \gamma_p \sin p f$$

where

$$\gamma_p = \frac{\overline{\cos p f}}{p}$$

$$\overline{\cos p f} = \left( \frac{-e}{1+\eta} \right)^p (1+p\eta)$$

$$\eta = \sqrt{1-e^2}$$

Appendix 2

Symbols For Functions

$$A_{\ell} = \left(\frac{a}{\lambda}\right)^{\ell+1} \cos(\beta + \theta \frac{f}{f_0})$$

$$A_n = \left(\frac{a}{\lambda}\right)^{n+1} \cos(\alpha + \theta \frac{f}{f_0})$$

$$\bar{A}_{\ell} = a_{\ell-1, \theta} \cos \beta$$

$$\bar{A}_n = a_{n-1, \theta} \cos \alpha$$

$$B_{KP} = f - M$$

$$\bar{B}_{KP} = 0$$

$$C_{\ell} = \frac{\cos[\beta + (\ell+k)\frac{f}{f_0}]}{\ell+k} + \frac{\cos[\beta + (\ell-k)\frac{f}{f_0}]}{\ell-k} \quad \text{for } \ell-k \geq 0$$

$$C_{\ell} = \frac{\cos(\beta + 2\ell \frac{f}{f_0})}{2\ell} \quad \text{for } \ell = 0$$

$$\bar{C}_{\ell} = (\gamma_{\ell+k} + \gamma_{\ell-k}) \cos \beta \quad \gamma_0 = 0$$

$$D'_{\ell} = \left(\frac{a}{\lambda}\right)^{\ell} \{ \sin[\beta + (\ell+1)\frac{f}{f_0}] - \sin[\beta + (\ell-1)\frac{f}{f_0}] \}$$

$$\bar{D}'_{\ell} = (a_{\ell-2, \theta+1} - a_{\ell-2, \theta-1}) \sin \beta$$

$$D'_{\ell} = \left(\frac{a}{\lambda}\right)^{\ell+1} \{ \sin[\beta + (\ell+1)\frac{f}{f_0}] - \sin[\beta + (\ell-1)\frac{f}{f_0}] \}$$

$$\bar{D}'_{\ell} = (a_{\ell-3, \theta+1} - a_{\ell-3, \theta-1}) \sin \beta \quad \text{for } \ell \geq 3$$

$$\bar{D}'_{\ell} = \left\{ \overline{\cos(\ell+1)\frac{f}{f_0}} - \overline{\cos(\ell-1)\frac{f}{f_0}} + \frac{a}{2} [\overline{\cos(\ell+2)\frac{f}{f_0}} - \overline{\cos(\ell-2)\frac{f}{f_0}}] \right\} \frac{\sin \beta}{\eta^2} \quad \text{for } \ell=2$$

$$E = -\frac{(n+1)}{2} \frac{a}{\eta} \left(\frac{a}{\lambda}\right)^{n+2} \cos[\alpha + (\theta+1)\frac{f}{f_0}] + \frac{(n+1)}{2} \frac{a}{\eta} \left(\frac{a}{\lambda}\right)^{n+2} \cos[\alpha + (\theta-1)\frac{f}{f_0}] - \frac{a}{\eta} \left(\frac{a}{\lambda}\right)^{n+2} \cos(\alpha + \theta \frac{f}{f_0})$$

$$\bar{E} = \left\{ -\frac{(n+1)}{2} \frac{a}{\eta} a_{n, \theta+1} + \frac{(n+1)}{2} \frac{a}{\eta} a_{n, \theta-1} - \frac{a}{\eta} a_{n, \theta} \right\} \cos \alpha$$



$$F = \left\{ \left( \frac{n+1+q}{2} \right) \left( \frac{a}{n} \right)^{n+2} + \frac{q}{2\eta^2} \left( \frac{a}{n} \right)^{n+1} \right\} \cos [\alpha + (q+1)f] \\ + \left\{ \left( \frac{n+1-q}{2} \right) \left( \frac{a}{n} \right)^{n+2} - \frac{q}{2\eta^2} \left( \frac{a}{n} \right)^{n+1} \right\} \cos [\alpha + (q-1)f]$$

$$\bar{F} = \left\{ \left( \frac{n+1+q}{2} \right) a_{n,q+1} + \frac{q}{2\eta^2} a_{n-1,q+1} + \left( \frac{n+1-q}{2} \right) a_{n,q-1} - \frac{q}{2\eta^2} a_{n-1,q-1} \right\} \cos \alpha$$

$$G' = \left\{ -\frac{e(n+1+q)(n+1-q)}{4\eta^2} \left( \frac{a}{n} \right)^{n+3} - \frac{e(n+1)q}{4\eta^2} \left( \frac{a}{n} \right)^{n+2} \right\} \sin [\alpha + (q+2)f] \\ + \left\{ -\frac{e(n+1+q)(n+1-q)}{4\eta^2} \left( \frac{a}{n} \right)^{n+3} - \frac{e(n+1)q}{4\eta^2} \left( \frac{a}{n} \right)^{n+2} \right\} \sin [\alpha + (q+1)f] \\ + \left\{ \frac{e(n+1+q)(n+1-q)}{4\eta^2} \left( \frac{a}{n} \right)^{n+3} + \frac{e(n+1)q}{4\eta^2} \left( \frac{a}{n} \right)^{n+2} \right\} \sin [\alpha + qf] \\ + \left\{ -\frac{e(n+1+q)(n+1-q)}{4\eta^2} \left( \frac{a}{n} \right)^{n+3} + \frac{e(n+1)q}{4\eta^2} \left( \frac{a}{n} \right)^{n+2} \right\} \sin [\alpha + (q-1)f] \\ + \left\{ \frac{e(n+1+q)(n+1-q)}{4\eta^2} \left( \frac{a}{n} \right)^{n+3} - \frac{e(n+1)q}{4\eta^2} \left( \frac{a}{n} \right)^{n+2} \right\} \sin [\alpha + (q-2)f]$$

$$\bar{G}' = 0$$

$$H = \left\{ \frac{(n+q+1)(n+q+3)}{4} \left( \frac{a}{n} \right)^{n+3} + \frac{(2nq+2q^2+4q+2n+1)}{4\eta^2} \left( \frac{a}{n} \right)^{n+2} + \frac{q(q+1)}{4\eta^2} \left( \frac{a}{n} \right)^{n+1} \right\} \cos [\alpha + (q+2)f] \\ + \left\{ \frac{(n+q+1)(n+q+3)}{4} \left( \frac{a}{n} \right)^{n+3} - \frac{(2nq+2q^2+4q+2n+1)}{4\eta^2} \left( \frac{a}{n} \right)^{n+2} + \frac{q(q+1)}{4\eta^2} \left( \frac{a}{n} \right)^{n+1} \right\} \cos [\alpha + (q-2)f] \\ + \left\{ \frac{(n+q+1)(n+q+3)}{4} \left( \frac{a}{n} \right)^{n+3} - \frac{(q^2+n+1)}{2\eta^2} \left( \frac{a}{n} \right)^{n+2} - \frac{q^2}{2\eta^2} \left( \frac{a}{n} \right)^{n+1} \right\} \cos (\alpha + qf)$$

$$\bar{H} = - \left\{ \frac{b_{n-1,q}''}{\eta} - \frac{(2n+1)}{\eta} \left[ 2 \frac{b_{n-1,q}}{\eta} + \frac{a_{n-1,q}}{n} \left\{ \frac{e}{2} - (2n+1) \right\} \right] \right\} \cos \alpha$$

$$W = \left[ (2\ell-1) B_{\ell,s} - \frac{\cos^2 \alpha}{\sin i} B'_{\ell,s} \right] a'_{\ell-1,s} + \frac{q}{e} \ell'_{\ell-1,s} B_{\ell,s}$$

Appendix 3

Derivatives of the Disturbing Function

$$R_n = \frac{\mu J_n}{a^{n+1}} \sum_q B_{n,q} (A - \bar{A})$$

$$\frac{\partial R_n}{\partial a} = -\frac{\mu(n+1)J_n}{a^{n+2}} \sum_q B_{n,q} (A - \bar{A})$$

$$\frac{\partial R_n}{\partial e} = \frac{\mu J_n}{a^{n+1}} \sum_q B_{n,q} (F - \bar{F})$$

$$\frac{\partial R_n}{\partial i} = \frac{\mu J_n \cos i}{a^{n+1}} \sum_q B'_{n,q} (A - \bar{A})$$

$$\frac{\partial R_n}{\partial \omega} = -\frac{\mu J_n}{a^{n+1}} \sum_q q B_{n,q} (A' - \bar{A}')$$

$$\frac{\partial R_n}{\partial M} = \frac{\mu J_n}{a^{n+1}} \sum_q B_{n,q} E'$$

$$\frac{\partial^2 R_n}{\partial a^2} = \frac{\mu(n+1)(n+2)J_n}{a^{n+3}} \sum_q B_{n,q} (A - \bar{A})$$

$$\frac{\partial^2 R_n}{\partial a \partial e} = -\frac{\mu(n+1)J_n}{a^{n+2}} \sum_q B_{n,q} (F - \bar{F})$$

$$\frac{\partial^2 R_n}{\partial a \partial i} = -\frac{\mu(n+1)J_n \cos i}{a^{n+2}} \sum_q B'_{n,q} (A - \bar{A})$$

$$\frac{\partial^2 R_n}{\partial a \partial \omega} = \frac{\mu(n+1)J_n}{a^{n+2}} \sum_q q B_{n,q} (A' - \bar{A}')$$

$$\frac{\partial^2 R_n}{\partial a \partial M} = -\frac{\mu(n+1)J_n}{a^{n+2}} \sum_q B_{n,q} E'$$

$$\frac{\partial^2 R_n}{\partial e \partial a} = -\frac{\mu(n+1)J_n}{a^{n+2}} \sum_q B_{n,q} (F - \bar{F})$$

$$\frac{\partial^2 R_n}{\partial e^2} = \frac{\mu J_n}{a^{n+1}} \sum_q B_{n,q} (H - \bar{H})$$

$$\frac{\partial^2 R_n}{\partial e \partial e} = \frac{\mu J_n \cos i}{a^{n+1}} \sum_q B'_{n,q} (F - \bar{F})$$

$$\frac{\partial^2 R_n}{\partial e \partial \omega} = -\frac{\mu J_n}{a^{n+1}} \sum_q q B_{n,q} (F' - \bar{F}')$$

$$\frac{\partial^2 R_n}{\partial e \partial H} = \frac{\mu J_n}{a^{n+1}} \sum_q B_{n,q} G'$$

$$\frac{\partial^2 R_n}{\partial a \partial a} = -\frac{\mu(n+1)J_n \cos i}{a^{n+2}} \sum_q B'_{n,q} (A - \bar{A})$$

$$\frac{\partial^2 R_n}{\partial a \partial e} = \frac{\mu J_n \cos i}{a^{n+1}} \sum_q B'_{n,q} (F - \bar{F})$$

$$\frac{\partial^2 R_n}{\partial a \partial \omega} = -\frac{\mu J_n \cos i}{a^{n+1}} \sum_q B'_{n,q} (A' - \bar{A}')$$

$$\frac{\partial^2 R_n}{\partial a \partial H} = \frac{\mu J_n \cos i}{a^{n+1}} \sum_q B'_{n,q} E'$$

$$\frac{\partial^2 R_n}{\partial \omega \partial a} = \frac{\mu(n+1)J_n}{a^{n+2}} \sum_q q B_{n,q} (A' - \bar{A}')$$

$$\frac{\partial^2 R_n}{\partial \omega \partial e} = -\frac{\mu J_n}{a^{n+1}} \sum_q q B_{n,q} (F - \bar{F})$$

$$\frac{\partial^2 R_n}{\partial \omega \partial a} = -\frac{\mu J_n \cos i}{a^{n+1}} \sum_q q B'_{n,q} (A' - \bar{A}')$$

$$\frac{\partial^2 R_n}{\partial \omega^2} = -\frac{\mu J_n}{a^{n+1}} \sum_q q^2 B_{n,q} (A - \bar{A})$$

$$\frac{\partial^2 R_n}{\partial \omega \partial H} = \frac{\mu J_n}{a^{n+1}} \sum_q q B_{n,q} E$$

# Appendix 4

## Perturbations in the Elements

$$\frac{\delta a}{a} = \frac{2J_2}{a^4} \sum_s B_{s,0} (A_s - \bar{A}_s)$$

$$\delta e = (\delta e)_1 + (\delta e)_2 + (\delta e)_3$$

$$(\delta e)_1 = \frac{\eta^2 J_2}{2 a^4} \sum_s B_{s,0} (A_s - \bar{A}_s)$$

$$(\delta e)_2 = \frac{J_2}{e a^2} \sum_s s a_{s-1,s} B_{s,0} B_{kp} \sin \beta$$

$$(\delta e)_3 = -\frac{J_2}{e a^2} \sum_s \sum_{k=0}^{s-1} s a_{s-1,k} \epsilon_k B_{s,0} (C_k - \bar{C}_k)$$

$$\delta i = (\delta i)_1 + (\delta i)_2$$

$$(\delta i)_1 = -\frac{J_2 \cos i}{a^2 \eta \sin i} \sum_s s a_{s-1,s} B_{s,0} B_{kp} \sin \beta$$

$$(\delta i)_2 = \frac{J_2 \cos i}{a^2 \eta \sin i} \sum_s \sum_{k=0}^{s-1} s a_{s-1,k} \epsilon_k B_{s,0} (C_k - \bar{C}_k)$$

$$\delta \omega = (\delta \omega)_1 + (\delta \omega)_2 + (\delta \omega)_3 + (\delta \omega)_4$$

$$(\delta \omega)_1 = \frac{J_2}{\eta^2 a^4} \sum_s W B_{s,0} \cos \beta$$

$$(\delta \omega)_2 = \frac{J_2}{\eta^2 a^4} \sum_s \sum_{k=0}^{s-1} \epsilon_k W (C_k' - \bar{C}_k')$$

$$(\delta \omega)_3 = \frac{J_2}{2 e a^4} \sum_s B_{s,0} (\bar{D}_{1s}' - \bar{D}_{1s})$$

$$(\delta \omega)_4 = \frac{J_2}{2 e a^4} \sum_s B_{s,0} (\bar{D}_{2s}' - \bar{D}_{2s})$$



$$\delta H = (\delta H)_1 + (\delta H)_2 + (\delta H)_3 + (\delta H)_4$$

$$(\delta H)_1 = -\frac{J_e}{e\eta^{1/2}a^2} \sum_s t'_{s+1,s} B_{s,s} \cos \beta$$

$$(\delta H)_2 = -\frac{J_e}{e\eta^{1/2}a^2} \sum_s \sum_{k=0}^{s-1} \epsilon_k t'_{s+1,k} B_{s,s} (C'_s - \bar{C}'_s)$$

$$(\delta H)_3 = -\frac{\gamma J_e}{e a^2} \sum_s B_{s,s} (D'_{s,s} - \bar{D}'_{s,s})$$

$$(\delta H)_4 = -\frac{J_e}{e\eta a^2} \sum_s B_{s,s} (D'_{s,s} - \bar{D}'_{s,s})$$

Appendix 5

Some Useful Trigonometric Identities

$$f - H = -2 \sum_{t=1}^{\infty} r_t \sin t f$$

$$\cos(\alpha + q f) \sin t f = \frac{1}{2} \{ \sin[\alpha + (q+t)f] - \sin[\alpha + (q-t)f] \}$$

$$(f - H) \cos(\alpha + q f) = \sum_{t=1}^{\infty} r_t \{ \sin[\alpha + (q+t)f] - \sin[\alpha + (q-t)f] \}$$

$$(f - H) \sin(\alpha + q f) = \sum_{t=1}^{\infty} r_t \{ \cos[\alpha + (q+t)f] - \cos[\alpha + (q-t)f] \}$$

$$\sin(\alpha + q f) \cos(\beta + s f) = \frac{\sin[\alpha + \beta + (q+s)f]}{2} + \frac{\sin[\alpha - \beta + (q-s)f]}{2}$$

$$\cos(\alpha + q f) \sin(\beta + s f) = \frac{\sin[\beta + \alpha + (q+s)f]}{2} + \frac{\sin[\beta - \alpha + (q-s)f]}{2}$$

$$\sin(\alpha + q f) \sin(\beta + s f) = \frac{\cos[\alpha - \beta + (q-s)f]}{2} - \frac{\cos[\alpha + \beta + (q+s)f]}{2}$$

$$\cos(\alpha + q f) \cos(\beta + s f) = \frac{\cos[\alpha - \beta + (q-s)f]}{2} + \frac{\cos[\alpha + \beta + (q+s)f]}{2}$$

# APPENDIX 6 NOMENCLATURE

Functions  $A_n, A_L, C_L, D_{1L}, D_{2L}, E_n, F_n, G_n, H_n$  are defined in Appendix 2. The primed function when not given explicitly is formed from the unprimed function. When the function has the subscript  $n$  then the primed function is formed by replacing  $\alpha$  in the unprimed function by  $\alpha - \frac{\pi}{2}$ . Similarly, when the function has the subscript  $l$  the primed function is formed by replacing  $\beta$  in the unprimed function by  $\beta - \frac{\pi}{2}$ .

- $B_{\epsilon p}$  = f-M = equation of the center
- $B_{n,q}$  = the inclination function
- $B'_{n,q}$  the derivative of  $B_{n,q}$  with respect to sine of  $i$ .
- $B''_{n,q}$  the second derivative of  $B_{n,q}$  with respect to sine of  $i$ .
- $J_n$  zonal harmonic of order  $n$
- $M$  mean anomaly
- ODL function which is zero for  $l$  even, and equals one for  $l$  odd.
- ODN function which is zero for  $n$  even, and equals one for  $n$  odd.
- $R$  disturbing function
- $R_L$  disturbing function of  $l^{th}$  zonal harmonic
- $R_n$  disturbing function of  $n^{th}$  zonal harmonic
- $W$  function defined in Appendix 2

$a$  semi major axis

$a'_{p-1,q}$  function defined in Appendix 1

$$a'_{p-1,q} = \frac{a^{p-1,q}}{\eta^{2p-4}}$$

$$b'_{p-1,q} = \frac{da_{p-1,q}}{de}$$

$$b''_{p-1,q} = \frac{d^2 a'_{p-1,q}}{de^2}$$

$e$  eccentricity

$i$  inclination

$i$  summation index ( when it cannot be confused with the inclination)

$j$  summation index

$$q = 2i + 0.0N, \quad i = 0, 1, 2, \dots, \frac{2-0.0N}{2}$$

$r$  radius vector of satellite

$$s = 2j + 0.0L, \quad j = 0, 1, 2, \dots, \frac{L-0.0L}{2}$$

$$\alpha = g\omega - 0.0N \frac{\pi}{2}$$

$$\beta = g\omega - 0.0L \frac{\pi}{2}$$

$\gamma_g$  coefficient appearing in equation of the center

$$\epsilon_g = \frac{1}{2} \text{ for } g=0$$

$$= 1 \text{ for } g>0$$

$$\eta = \sqrt{1-e^2}$$

$\chi(d), \phi(d), \rho(d), \theta(d), \sigma(d), \tau(d)$  subscripts appearing in averaging functions.

$\omega$  argument of perigee

$\Omega$  right ascension of node